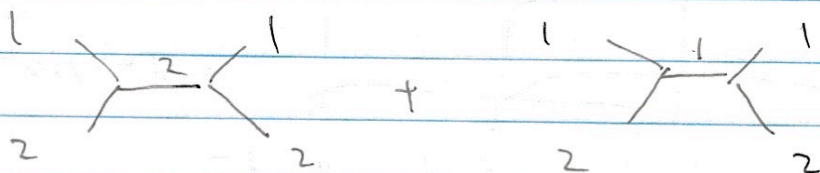


Schneitz
7.6
(in progress)

$$\mathcal{L}_{int} = \frac{\lambda}{2!} \phi_1 (\partial_\mu \phi_2) (\partial^\mu \phi_2) + \frac{g}{2} \phi_1^2 \phi_2$$

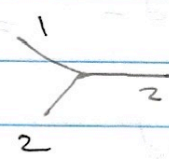
There are two types of diagrams, at tree level for $\phi_1 \phi_2 \rightarrow \phi_1 \phi_2$



$$\phi_2 = \int \frac{d^3k}{(2\pi)^3} \frac{1}{\sqrt{2\omega_k}} \left[a_k e^{ikx} + a_k^\dagger e^{-ikx} \right]$$

$$\partial_\mu \phi_2 = \int \frac{d^3k}{(2\pi)^3} \frac{1}{\sqrt{2\omega_k}} \left[ik_\mu a_k e^{ikx} - ik_\mu a_k^\dagger e^{-ikx} \right]$$

\Rightarrow annihilates ϕ_2 particle puts up ik_μ ,
creates ϕ_2 particle puts up $-ik_\mu$,

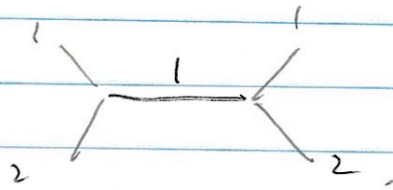
For our concern, we only need to consider the  vertex, where one ϕ_2 is annihilated,

and one ϕ_2 is created, so we get an overall factor

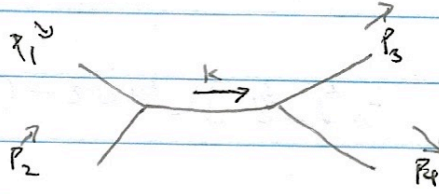
$$\left[ik_\mu (\text{incoming}) \right] \left[-ik_\mu (\text{outgoing}) \right]$$

$$= k_\mu (\text{incoming}) k^\mu (\text{outgoing})$$

Now for the diagram



In p -space:



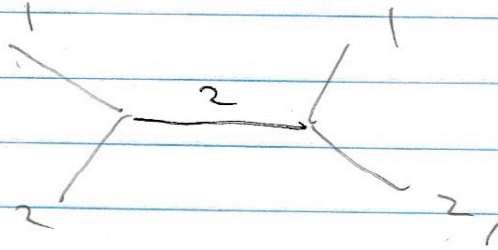
$$iM_2 = i g^2 \frac{i}{k^2 + i\epsilon} i g^2 \delta(p_1 + p_2 - k) \delta(k - p_3 - p_4)$$

$$M_2 = -g^2 \frac{1}{(p_1 + p_2)^2 + i\epsilon} \times \delta^4(\sum p)$$

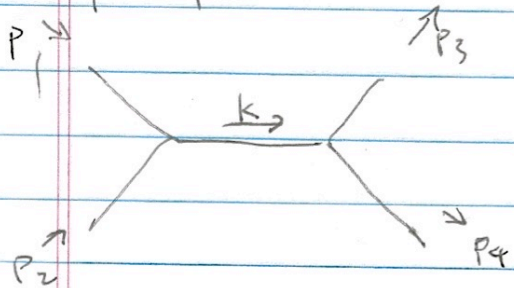
$$\Rightarrow M = M_1 + M_2$$

$$= \left[\frac{\lambda^2 p_2^\mu (p_1 + p_2)_\mu p_4^\nu (p_1 + p_2)_\nu}{(p_1 + p_2)^2 + i\epsilon} - g^2 \frac{1}{(p_1 + p_2)^2 + i\epsilon} \right] \times \delta^4(\sum p)$$

Now we evaluate



in p -space:



$$iM_1 = i\lambda (P_2^\mu k_\mu) \frac{i}{k^2 + i\epsilon} i\lambda (P_4^\nu k_\nu) \delta(P_1 + P_2 - k) \delta(k - P_3 - P_4)$$

$$= \frac{(-i\lambda^2) P_2^\mu (P_1 + P_2)_\mu P_4^\nu (P_1 + P_2)_\nu \times \delta(\Sigma p)}{(P_1 + P_2)^2 + i\epsilon}$$

$$M_1 = \frac{(-\lambda^2) P_2^\mu (P_1 + P_2)_\mu P_4^\nu (P_1 + P_2)_\nu \times \delta(\Sigma p)}{(P_1 + P_2)^2 + i\epsilon}$$

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